

AUG 22 2022

BULK & EDGE INDICES FOR 2D TOPOLOGICAL SYSTEMS

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QUANTISSIMA, VENICE 2022

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EXAMPLE $\frac{1}{2}$ (Bernoulli-Zhang-Hughes) Model

On $\ell^2(\mathbb{Z}^2)$, two right shift op. : e_1, e_2 std. basis of $\mathbb{R}^2$

$$(R_j \psi)(x) \equiv \psi(x - e_j) \quad (x \in \mathbb{Z}^2, j=1,2)$$

Define on $\ell^2(\mathbb{Z}^2) \otimes \mathbb{C}^2$, $a \in \mathbb{R}$, $\sigma_1, \sigma_2, \sigma_3$ Pauli mat.

$$\begin{aligned} H^{\frac{1}{2}BHZ} &:= a \mathbb{1} \otimes \sigma_3 + [R_1 \otimes \frac{1}{2}(\sigma_3 - i\sigma_1) + R_2 \otimes \frac{1}{2}(\sigma_3 - i\sigma_2) + \text{h.c.}] \\ &= (a \mathbb{1} + \text{Re}\{R_1 + R_2\}) \otimes \sigma_3 + \text{Im}\{R_1\} \otimes \sigma_1 + \text{Im}\{R_2\} \otimes \sigma_2 \end{aligned}$$

$$\text{Re}\{A\} \equiv \frac{1}{2}(A + A^*) \quad , \quad \text{Im}\{A\} \equiv \frac{1}{2i}(A - A^*)$$

FACT: If $a \neq 0, \pm 2$, $\exists$ gap for $H^{\frac{1}{2}BHZ}$ about zero energy.

If $\exists$ gap, define Hall conductivity

$$2\pi \sigma_{\text{Hall}} = 2\pi i \operatorname{tr} (P [\Lambda_1, P], [\Lambda_2, P])$$

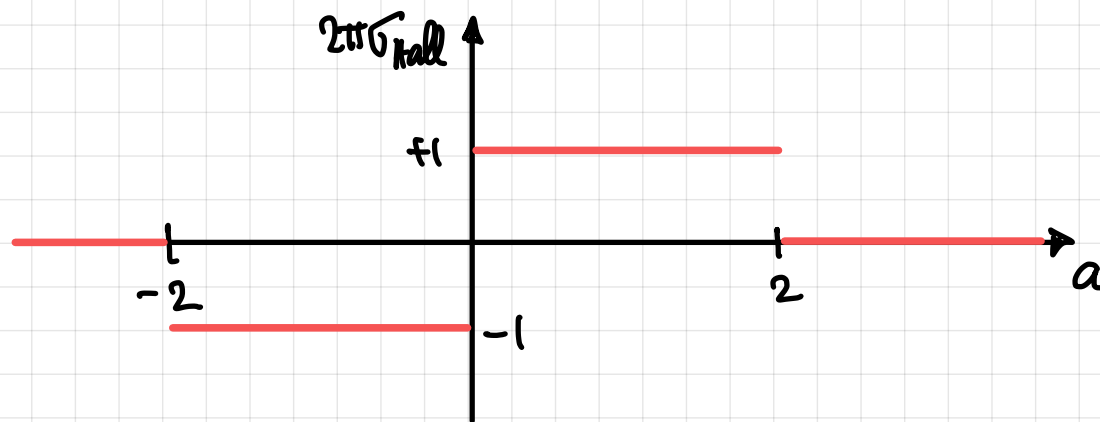
$$= \operatorname{index} (\underbrace{PUP + P}_{=: F}) \equiv \dim \ker F - \dim \ker F^* \in \mathbb{Z}$$

$$\omega / \quad P := \chi_{(-\infty, 0)}(H)$$

$$U \equiv \exp(i \arg(X_1 + iX_2))$$

$$\Lambda_j \equiv \chi_N(X_j) \quad j=1,2$$

FACT:



This is the simplest non-trivial IQHE tight-binding model.

Now make this model Time-Reversal-Invariant (TRI):

Full Bernevig-Zhang-Hughes Model

On $\ell^2(\mathbb{Z}^2) \otimes \mathbb{C}^2 \otimes \mathbb{C}^2$, define

$$H^{\text{BHZ}} := (a\mathbb{1} + \operatorname{Re}\{R_1 + R_2\}) \otimes \sigma_3 \otimes \mathbb{1} + \operatorname{Im}\{R_1\} \otimes \sigma_1 \otimes \sigma_3 + \operatorname{Im}\{R_2\} \otimes \sigma_2 \otimes \mathbb{1}$$

Note: H^{BHZ} is NOT block diagonal $\leadsto$ spin not conserved.

FACT: If $a \neq 0, \pm 2$, H^{BHZ} has a gap about zero.

Define Time-Reversal-Symmetry $\Theta := \underbrace{C}_{\substack{\uparrow \\ \text{complex-conj. of scalars}}} \mathbb{1} \otimes \mathbb{1} \otimes (-i\sigma_2)$

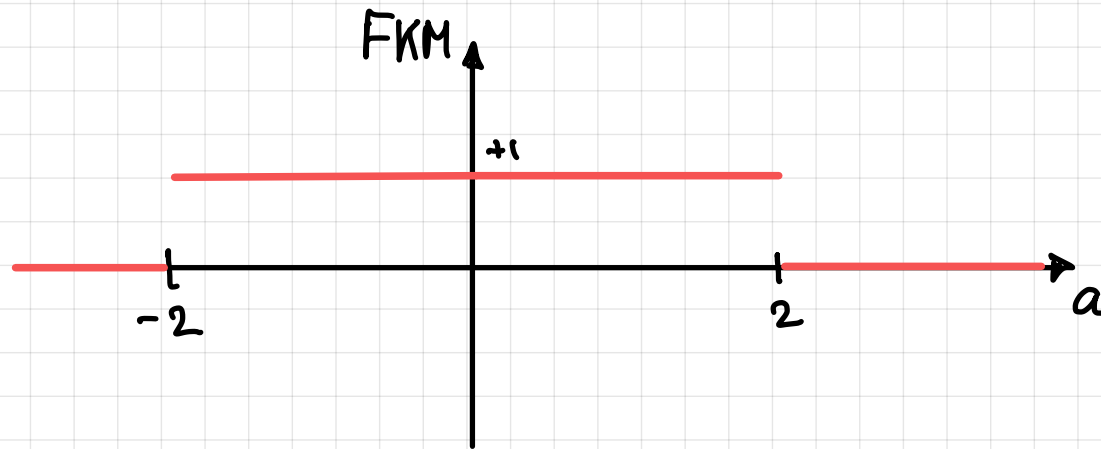
$$\Theta^2 = -\mathbb{1} \quad \leadsto \text{Fermionic TRS.}$$

FACT: $[H^{\text{BHZ}}, \Theta] = 0 \quad \leadsto \quad \sigma_{\text{Hall}}(H^{\text{BHZ}}) = 0 \quad \forall a.$

FU-KANE-MELE (2007) have defined another topological index in this context — $\mathbb{Z}_2$ -valued index counting whether edge states are unpaired.

$$FKM := \text{index}_2 \underbrace{(PUP + P^\perp)}_{\equiv F} \equiv ([\dim \ker F] \bmod 2) \in \mathbb{Z}_2$$

$$P \equiv \chi_{(-\infty, 0)}(H) \quad U \equiv \exp(i \arg(X_1 + iX_2))$$



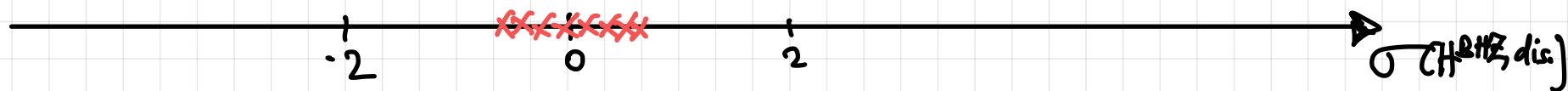
DISORDERED VERSION

Let $\{w(x)\}_{x \in \mathbb{Z}^2}$ be an IID seq. of R.V. w/ some nice prob. measure ρ .

$$H^{\text{BHZ, dis.}} := H^{\text{BHZ}} + w(X) \mathbb{1} \otimes \sigma_3 \otimes \mathbb{1}$$

Effect: make the parameter "a" vary randomly across sites.

Depending on ρ , $H^{\text{BHZ, dis.}}$ will exhibit Anderson localization about zero.



Results about general models:

Theorem(s): If H exhibits Anderson loc. about zero, $\mathbb{Z}$ and $\mathbb{Z}_2$ indices remain well-defined, obey a bulk-edge correspondence, and are independent of the choice of Fermi energy within the localized interval (=the mobility gap).

TECHNICAL

DETAILS

DEF.: $A \in \mathcal{B}(\ell^2(\mathbb{Z}^d))$ is called weakly-local iff $A \in WLOC$ iff

$$\exists \vartheta \in \mathbb{N}: \quad \forall \alpha \in \mathbb{N}, \exists C_\alpha \in (0, \infty):$$

$$\|A_{xy}\| \leq C_\alpha (1 + \|x - y\|)^{-\alpha} (1 + \|x\|)^{\vartheta}$$

For any interval $\Delta \subseteq \mathbb{R}$, let $B_1(\Delta)$ be the set of measurable f^n $f: \mathbb{R} \rightarrow \mathbb{C}$:

- ① f is const. above and below Δ .

$$\text{② } \|f\|_\infty \leq 1.$$

DEF.: $H = H^*$ has a deterministic mobility gap on Δ if

- ① $f(H) \in WLOC$ with estimates uniform as $f \in B_1(\Delta)$.
- ② $\sigma(H) \cap \Delta$ has finite degeneracy.

Note: $\chi_{(-\infty, 0)} \in B_1(\Delta) \Rightarrow P \equiv \chi_{(-\infty, 0)}(H) \in \mathcal{K}LOC$ when
 H has a det. mob. gap on
 some $\Delta \ni 0$.

Thm.: (Aizenman-Graf, Graf-Elgart-Schenker, ...)

$P \in \mathcal{K}LOC \Rightarrow F \equiv PUP + P^\perp$ is Fredholm.

Proof: $PUP + P^\perp$ is Fredholm iff $[P, U] \in \mathcal{C}pt$.

Indeed, then $PU^*P + P^\perp$ is the parametrix:

$$\begin{aligned} \mathbb{1} - (PUP + P^\perp)(PU^*P + P^\perp) &= P - PUPU^*P \\ &= P(UU^* - UPU^*)P = PUP^\perp U^*P \\ &= [P, U]P^\perp U^*P \end{aligned}$$

But $[P, f(X)] \in \mathcal{C}pt$ whenever $P \in \mathcal{K}LOC$ and

$$|f(x) - f(y)| \leq D \frac{\|x - y\|}{1 + \|x\|}.$$

SULE basis (Fimon et al...)

$\{\psi_n\}_n$ is a SULE basis for H on Δ

iff $\exists$ loc. centers $\{x_n\}_n \subseteq \mathbb{Z}^d$ and $\exists \epsilon \in \mathbb{N} : \forall \alpha > 0$

$$\exists C_\alpha \in (0, \infty) : \|\psi_n(x)\| \leq C_\alpha (1 + \|x - x_n\|)^{-\alpha} (1 + \|x_n\|)^2$$

Lemma: If H has a det. mob. gap on Δ then
 $\exists$ SULE basis.

STABILITY THEOREM

Let $\mu \in \Delta$ and $P_\mu := \chi_{(-\infty, \mu)}(H)$.

Thm.: If H has a det. mob. gap on Δ then

$$\Delta \ni \mu \mapsto \text{index}_{(2)}(P_\mu U P_\mu + P_\mu^\perp)$$

is constant.

Pf.: Let $\mu' > \mu$, $P := P_\mu$, $Q := P_{\mu'} - P$
 $EU := PUP + P^\perp$.

$$P_{\mu'} U = PUP + QUQ + (P+Q)^\perp + PUQ + QUP$$

But PUQ is opt.. Indeed:

$$|PUQ|^2 \equiv QU^*PUQ = Q(\underbrace{U^*PU - P}_{\in \mathcal{J}_3})Q$$

$$\Rightarrow \text{index}_{(2)} P_U U = \text{index}_{(2)} PUP + QUQ + (P+Q)^\perp$$

$$\begin{aligned} \ker(PUP + QUQ + (P+Q)^\perp) &\cong (\ker PUP) \oplus (\ker QUQ) \\ &\cong (\ker P_U) \oplus (\ker QU) \end{aligned}$$

$$\text{But } \text{index}_{(2)} A \oplus B = \text{index}_{(2)} A + \text{index}_{(2)} B \pmod{2}.$$

Hence W.T.S. $\text{index}_{(2)} QU = 0$ whenever Q projects within the mobility gap.

Is $QU - \mathbb{1}$ compact? NO.

By the above, $\exists$ SULE basis for $\text{im}(Q)$.

Define $V: \text{im}(Q) \rightarrow \text{im}(Q)$

$$\psi_n \mapsto \exp(i \arg((x_n)_1 + i(x_n)_2))$$

localization ctr.

V is unitary on $\text{im}(Q) \Rightarrow \text{index}_{(2)} (QVQ + Q^\perp) = 0$.

Claim: $Q(U-V)Q \equiv Q(\exp(i \arg(X_1 + iX_2)) - \exp(i \arg((x_n)_1 + i(x_n)_2)))Q$

is cpt.

Proof: ψ_n decays very quickly away from x_n anyway...

But index_{x_2} is stable under cpt. perturbations.

OPEN PROBLEM

- ① Define a new metric d on $\mathcal{B}(\ell^2(\mathbb{Z}^2) \otimes \mathbb{C}^2)$ s.t.
- $$\mathcal{S} := \left\{ H = H^* \in \mathcal{B}(\ell^2(\mathbb{Z}^2) \otimes \mathbb{C}^N) \mid PUP + PL \text{ is Fredholm} \right\}$$

is an open set.

- ② Show that set of path-connected components $\pi_0(\mathcal{S})$ induced by d is bijective w/

$\mathbb{Z}$

IQHE

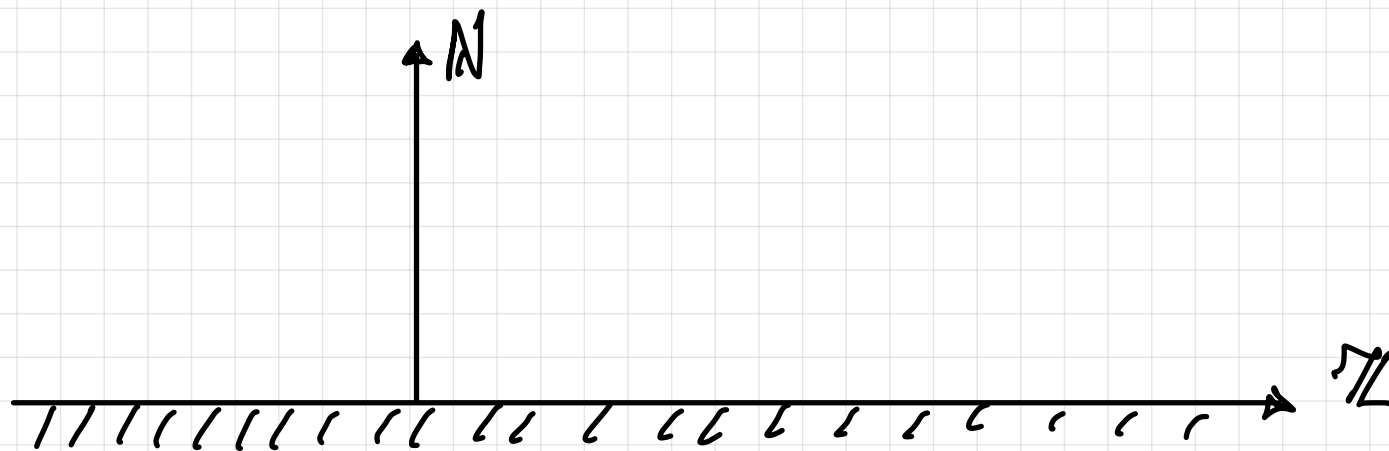
$\mathbb{Z}$

TRI

via $\text{index}_{(2)}$.

- ③ Relation of d to transport coeff. / moments of X ?

EDGE SYSTEMS



VACUUM

$$\hat{\mathcal{H}} := \ell^2(\mathbb{Z} \times N) \otimes \mathbb{C}^N$$

$$\iota : \ell^2(\mathbb{Z} \times N) \hookrightarrow \ell^2(\mathbb{Z}^2)$$

$\hookrightarrow$ injection

(extend by zero).

If H acts on $\mathcal{H}$, $\iota^* H \iota$ acts on $\hat{\mathcal{H}}$
(Dirichlet restriction).

As a rule, if H has a (spectral / mobility) gap, $\tau^* H \tau$ does NOT.

However, $\tau^* f(H) \tau - f(\tau^* H \tau) \rightarrow 0$ (into bulk)
for any smooth f .

$\Rightarrow$ If $\text{supp}(f) \subseteq \Delta$, then $f(\tau^* H \tau)$ decays into bulk.

Def.: $\hat{H} = \hat{H}^* \in \mathcal{B}(\mathcal{H})$ has a bulk-spec.-gap on Δ iff $(g^2 - g)(\hat{H})$ decays into bulk where $g \approx$ smooth version of $\chi_{(-\infty, 0]}$ with $\text{supp}(g^2 - g) \subseteq \Delta$.

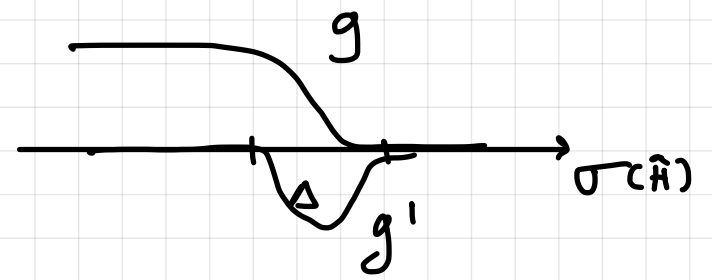
Unfortunately for the mob. gap regime this cannot work, since then, H has eigenvalues in Δ which are centered everywhere, so $(g^2 - g)(H) \neq 0$ and does not even decay into bulk.

Intrinsic def. for the edge???

DEF.: $\hat{H} = \hat{H}^* \in \mathcal{B}(\hat{\mathcal{H}})$ has a bulk-mob.-gap on Δ iff $\exists$ bulk Hamiltonian $K \in \mathcal{B}(\mathcal{H})$ w/ mob. gap on Δ such that $\hat{H} - \mathcal{I}^* K \mathcal{I} \rightarrow 0$ (into bulk).

This implies first def.

EDGE INDEX



QHE: Edge hall conductivity

$$2\pi \hat{\sigma}_{\text{Hall}} = 2\pi \text{tr} \left(g'(\hat{H}) \underbrace{i[\Lambda_1, \hat{H}]}_{\text{"velocity } v."} \right)$$

$$= \text{index} \left(\Lambda_1 \exp(-2\pi i g(\hat{H})) \Lambda_1 + \Lambda_1 \right) \in \mathbb{Z}$$

FKM: ~~trace~~ trace formula.

Define $\text{index}_{\text{FKM}}(\Lambda_1 \exp(-2\pi i g(\hat{H})) \Lambda_1 + \Lambda_1)$

Relies on $[\Lambda_1, 1 - e^{-2\pi i g(\hat{H})}] \in \text{cpt.}$

$\Leftarrow (g^2 - g)(\hat{H})$ decays in 2 direction

NOT true in mob. gap regime.

Idea: Remove localized bulk states from $g(\hat{H})$:

Since $\exists K$ w/ mob. gap on Δ :

$$\hat{H} - \mathbb{1}^* K \mathbb{1} \rightarrow 0 \quad \text{into the bulk}$$

$$\hat{R} := \mathbb{1}^* \chi_{\Delta^c}(K) \mathbb{1}$$

In spec. gap regime, $\hat{R} = \mathbb{1}$ as $\chi_{\Delta^c}(K) = \mathbb{1}$.

Consider $\hat{F} := \Lambda_1 e^{-2\pi i \hat{R} g(\hat{H}) \hat{R}} \Lambda_1 + \Lambda_1^\perp$.

Thm. 0 $\hat{F}$ is Fredholm and its index is indep. of K
as long as $\hat{H} - \mathbb{1}^* K \mathbb{1} \rightarrow 0$ into bulk.

Thm. 0 $\text{index}_{(2)} F = \text{index}_{(2)} \hat{F}$ if $\hat{H} - \mathbb{1}^* H \mathbb{1} \rightarrow 0$.