

Self-avoiding walk and the renormalisation group

David C. Brydges and Gordon Slade

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Constructive and Multiscale Methods in Quantum Theory

Self-avoiding walks on \mathbb{Z}^d are simple-random walk paths without self-intersections. Self-avoiding walks of the same length are declared to be equally likely. The basic question is how far on average is their endpoint from the origin?

After a brief review of the current state of knowledge I will describe work in progress with Gordon Slade for the case $d = 4$ which is an application of the renormalisation group to a supersymmetric lattice field theory. Our immediate goal is to prove that the critical two-point function (Green function) for a spread-out model of self-avoiding walks on \mathbb{Z}^d decays like $|x|^{-2}$ at large distances, as it does for simple random walk.

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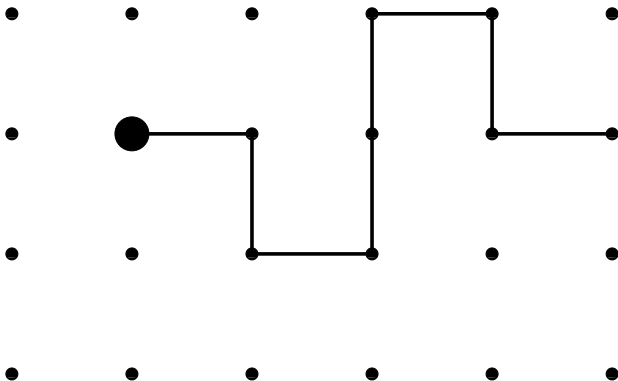


Figure: An $N = 8$ step self-avoiding walk

Equal probability for all N step self-avoiding walks.

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Defn: $a_N \sim b_N$ if $\lim_{N \rightarrow \infty} \frac{a_N}{b_N} \in (0, \infty)$.

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Brydges-Imbrie 2003: true for weak coupling Edwards model on **hierarchical** four dimensional lattice

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- ▶ At $\lambda = \lambda_c$, $G_{a,b} \sim |a - b|^{-(d-2)}$

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- ▶ For $g \geq 0$ there exists $\lambda_c = \lambda_c(g)$ s.t. divergent for $\lambda < \lambda_c$, convergent for $\lambda > \lambda_c$

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- ▶ The last step is an abelian version of

$$\mathbb{E}(\text{end-to-end distance}) \sim N^{1/2} \log^{1/8}(N)$$

Related results

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- ▶ D. Iagolnitzer and J. Magnen 1994: Continuum Edwards Model:

$$\text{Gr}_{ab}(z) \sim |a - b|^{-2} \left(1 + \lambda_1 \frac{1}{\ln |a - b|} + \lambda_2 \frac{\ln \ln |a - b|}{\ln^2 |a - b|} + r(a, b) \right)$$

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- ▶ Edwards model $d = 4$ hierarchical lattice: Suemi Rodríguez-Romo (1995); Hara, Ohno (2009).
- ▶ Edwards model based on Lévy walk on \mathbb{Z}^3 (model has upper critical dimension $3 + \epsilon$): Mitter, Scoppola (2008).

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- ▶ $V_x =$ replace ϕ by ϕ_x .
- ▶ $V(\Lambda) = \sum_{x \in \Lambda} V_x$.
- ▶ There exist C and V such that

$$G_{a,b} = \lim_{\Lambda \uparrow \mathbb{Z}^4} \mathbb{E} \left(e^{-V(\Lambda)} \bar{\varphi}_a \varphi_b \right)$$

References

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- ▶ Parisi-Sourlas (1980); McKane, (1980); Luttinger (1983);

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- ▶ Review: Brydges, Imbrie, Slade (2009).

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Brydges, Guadagni, Mitter (2004): existence for periodic lattice
 $C = (\epsilon - \Delta)^{-1}$ for $\epsilon > 0$.

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- ▶ where, on the right hand side,

$$\phi = \xi_1 + \xi_2 + \cdots + \xi_N$$

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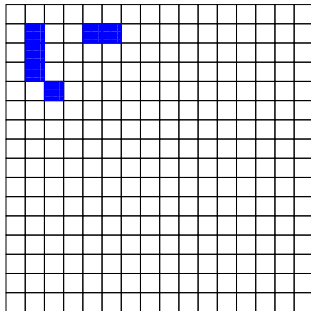
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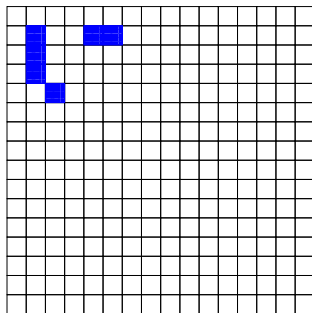
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- ▶ $Z_N = \mathbb{E} Z_0$ is our Green's function $G_{a,b}^\Lambda$

Lattice Geometry

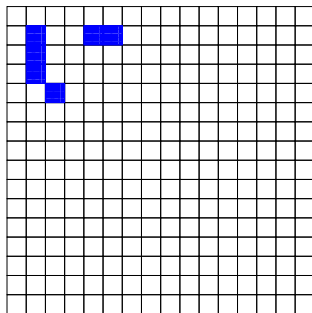


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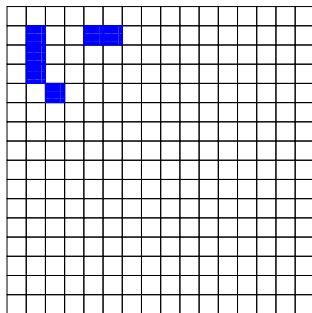
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- ▶ Let $\mathcal{P}_j = \{\text{finite unions of blocks in } \mathcal{B}_j\}$.
- ▶ Given functions I, K defined on \mathcal{P}_j , let

$$(I \circ K)(\Lambda) = \sum_{X \in \mathcal{P}_j} I(\Lambda \setminus X) K(X)$$

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 - ▶ $K_0(X) = \mathbb{1}_{X=\emptyset}$
 - ▶ $I_0(X) = e^{-V(X)}$
- ▶ What about $\bar{\varphi}_a \varphi_b$? Add $\sigma_a \bar{\varphi}_a + \bar{\sigma}_b \varphi_b$ to $V_0(\Lambda)$ so that

$$G_{a,b} = \left. \frac{\partial}{\partial \sigma_a} \frac{\partial}{\partial \bar{\sigma}_b} \right|_{\sigma=0} \mathbb{E} I_0 \circ K_0(\Lambda)$$

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- ▶ Given (I_j, K_j) there exist (I_{j+1}, K_{j+1}) such that

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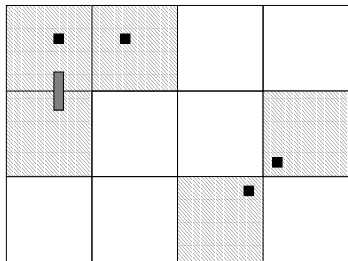
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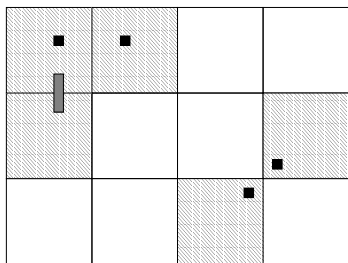
- ▶ $V_j \equiv (g_j, a_j, z_j, \sigma_j)$.

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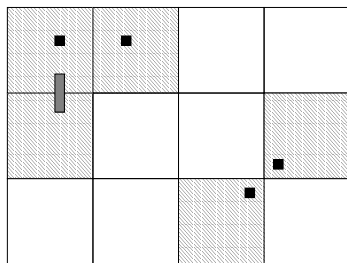


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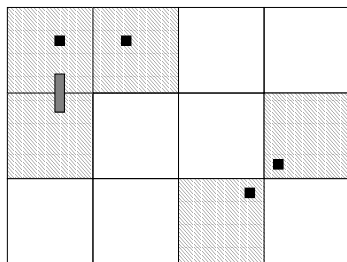
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- ▶ with \tilde{K} given by: for $U \in \mathcal{P}_{j+1}$,

$$\tilde{K}(U) = \sum_{X \in \overline{\mathcal{P}}_j(U)} I_{j+1}(U \setminus X) \mathbb{E}_{j+1}(K_j \circ \delta I)(X).$$

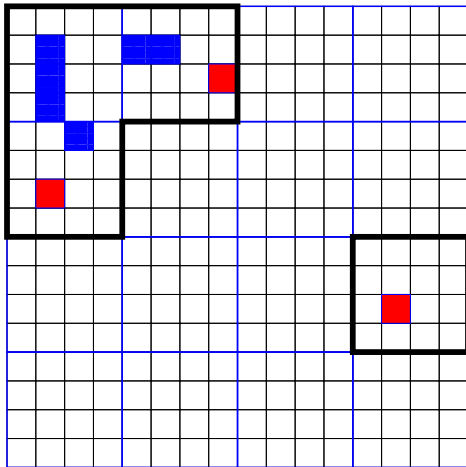
Picture illustrating $\tilde{K}(U)$

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Picture illustrating $\tilde{K}(U)$

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- ▶ Summary: **Factorisation** of K_j at scale $j \Rightarrow$ Factorisation of \tilde{K} at scale $j + 1$

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- ▶ Example: $\sum_{X: X \supset B} \|K(X)\| \ll \|V(B)\|$ for all $B \in \mathcal{B}_j$.
- ▶ Such estimates cannot survive an induction $j \rightarrow j+1$ unless V_{j+1} is chosen well

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- ▶ $z_{j+1} = \sum_{k \geq j} r_{z,k}(\xi)$ is summable

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