

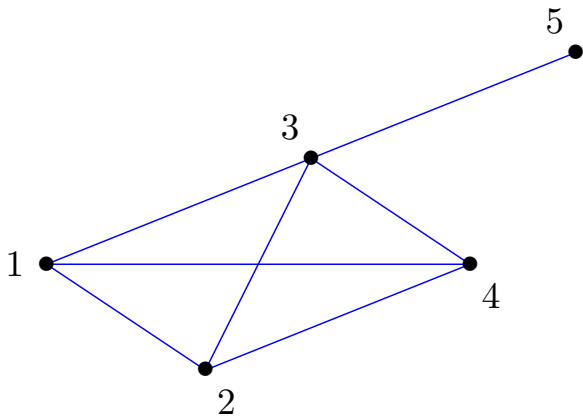
The random interchange process on the hypercube

Roman Kotecký

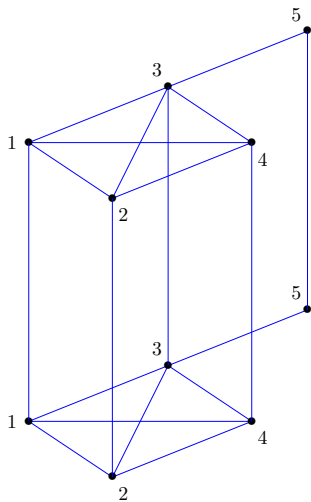
A joint work with P. Miłoś (Warsaw) and D. Ueltschi (Warwick)

Warwick/Prague

Random transpositions, resulting permutations and their cycles

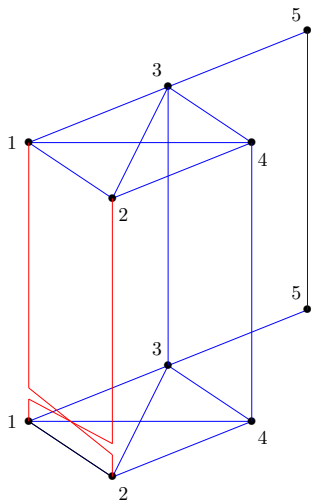


Random transpositions, resulting permutations and their cycles



$(1, 2, 3, 4, 5)$ $((1), (2), (3), (4), (5))$

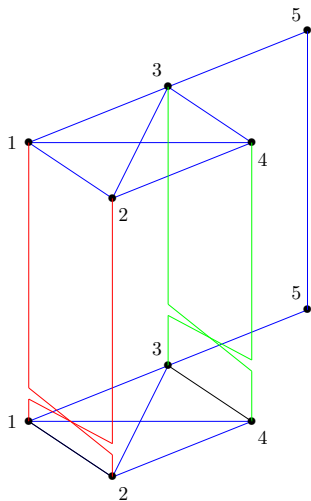
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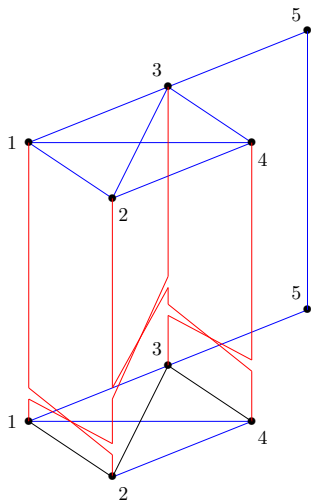


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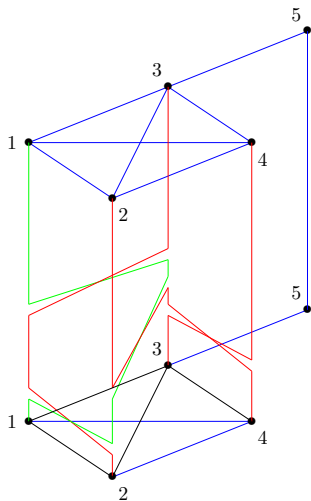
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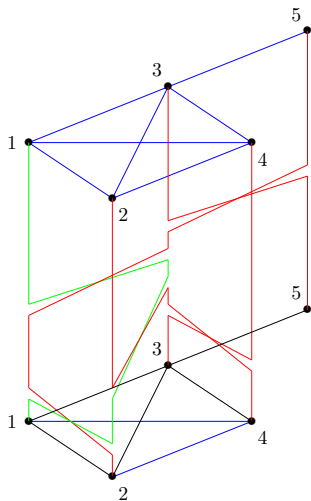
$(3, 1, 4, 2, 5)$ $((1, 3, 4, 2), (5))$

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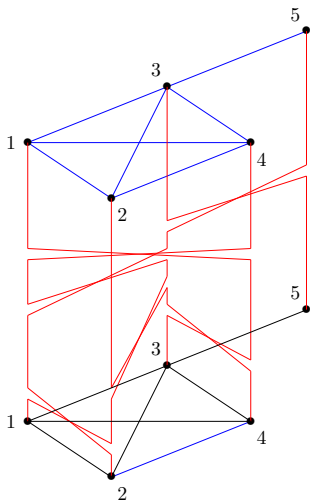
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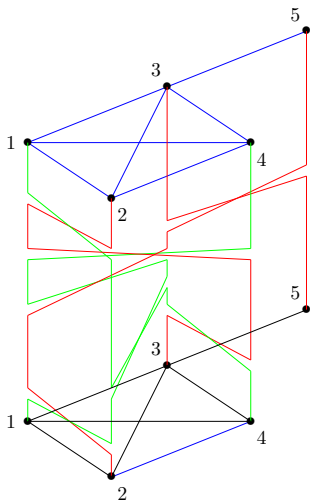
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Random transpositions on a complete graph

Card shuffling

$$K_N : \pi_t = T_t \circ T_{t-1} \circ \cdots \circ T_1.$$

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Diaconis, Shahshahani, 81

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Problem makes sense for an arbitrary graph: for large times t , do cycles of length $\sim N$ exist?

Connection with physics

Bálint Tóth, 93 : Heisenberg ferromagnet via random interchange process on \mathbb{Z}^d :

the spontaneous magnetization

$$m(\beta) = \frac{1}{2} \lim_{k \rightarrow \infty} \lim_{\Lambda \nearrow \mathbb{Z}^d} \frac{\mathbb{E}(\mathbb{1}_{\ell_\beta(0) > k} 2^{\#\text{of cycles}})}{\mathbb{E}(2^{\#\text{of cycles}})} \sim \text{density of hard-core Bose gas} - \frac{1}{2}$$

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Phase transition is linked with the occurrence of long cycles

Hypercube $Q_d = \{0, 1\}^d$, $N = 2^d$, $E = \frac{1}{2}dN$

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Subcritical t

Ajtai, Komlós, Szemerédi 82

Bollobás, Kohayakawa, Łuczak 92

HS 05, 06

BCHSS 06, HN 12

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No giant component \implies only small cycles

$$t = cN$$

$$p = \frac{cN}{Nd/2} = \frac{2c}{d} \implies t_c \sim \frac{1}{2}N$$

Theorem

Let $c < 1/2$ and $\epsilon > 0$. Then there exists d_0 such that

$$\mathbb{P}(|V_t(\kappa d)| = 0) > 1 - \epsilon \kappa^{-3/2}$$

for all $t \leq cN$, all $\kappa \geq \frac{2 \ln 2}{(1-2c)^2}$, and all $d > d_0$.

Supercritical t

For $t > \frac{1}{2}N$, there are "big" cycles in average for a small interval around t

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Let $c > \frac{1}{2}$ and let (Δ_d) be a sequence of positive numbers such that $\Delta_d d / \log d \rightarrow \infty$ as $d \rightarrow \infty$. Then there exist $\eta(c) > 0$ and d_0 such that for all $d > d_0$, all $T > cN$, and all $a > 0$, we have

$$\frac{1}{\Delta_d T} \sum_{t=T+1}^{\lfloor (1+\Delta_d)T \rfloor} \mathbb{E} \left(\frac{|V_t(N^a)|}{N} \right) \geq \eta(c) - a.$$

For $c > 1$, we can take $\eta(c) = \frac{1}{2}(1 - \frac{1}{c})$.

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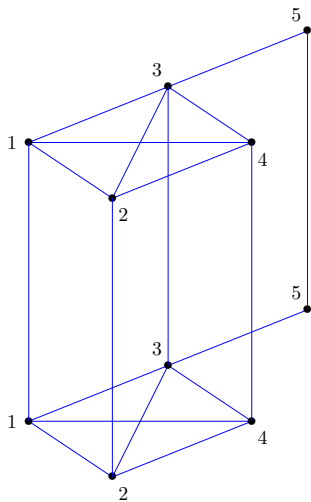
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The main obstacle:

Unlike percolation clusters, the cycles do not grow monotonously — an interchange over an edge with endpoints in a cycle splits it into two shorter cycles.

Dynamics of merging and splitting cycles

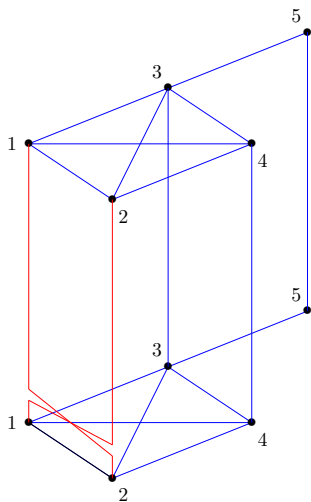


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$$N_t = \tilde{N}_t = 5$$

Dynamics of merging and splitting cycles

Merging clusters (and corresponding cycles): \widetilde{M}_t

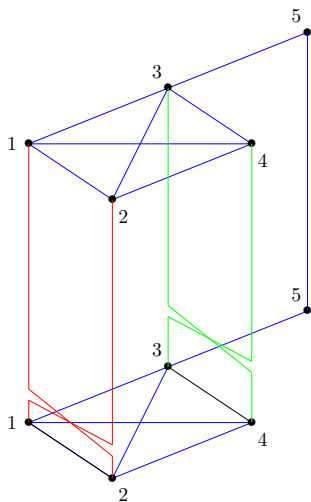


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$$N_t = \widetilde{N}_t = 4, \Delta N_t = -1, \Delta \widetilde{N}_t = -1$$

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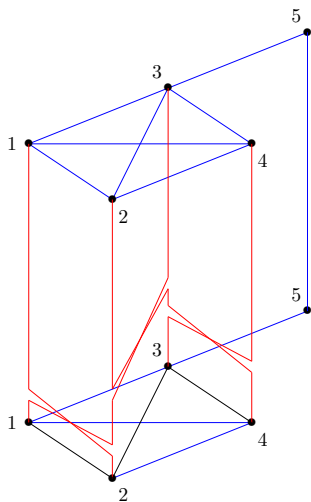


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$$N_t = \widetilde{N}_t = 3, \Delta N_t = -1, \Delta \widetilde{N}_t = -1$$

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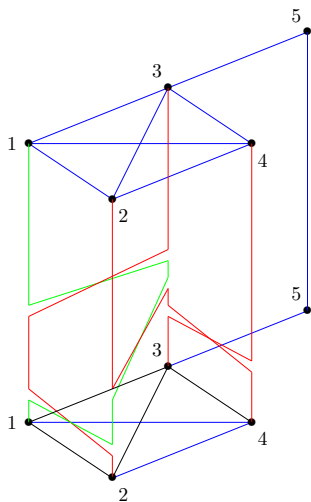


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$$N_t = \widetilde{N}_t = 2, \Delta N_t = -1, \Delta \widetilde{N}_t = -1$$

Dynamics of merging and splitting cycles

Splitting a cycle (within a cluster) S_t

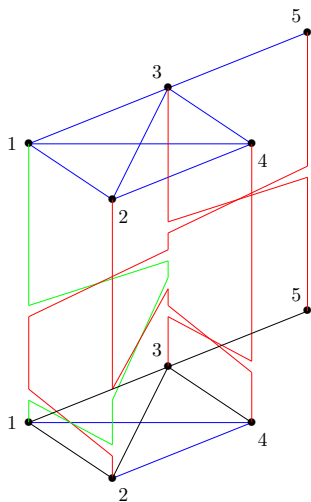


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$$N_t = 3, \tilde{N}_t = 2, \Delta N_t = 1, \Delta \tilde{N}_t = 0$$

Dynamics of merging and splitting cycles

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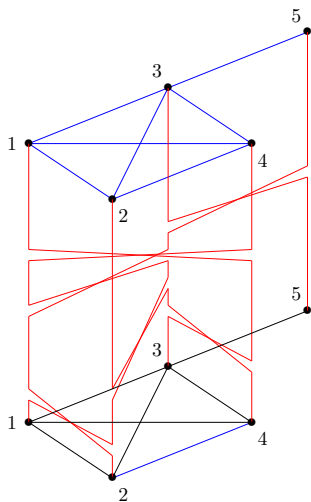


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Dynamics of merging and splitting cycles

Merging cycles (within a cluster) M_t

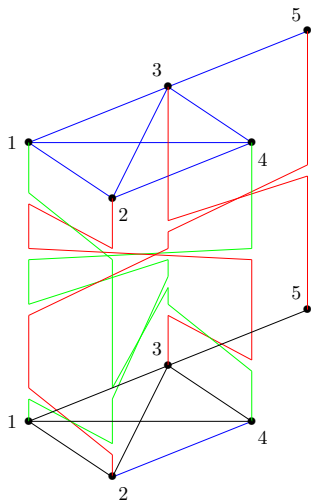


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Dynamics of merging and splitting cycles

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Supercritical t

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6 steps of the proof:

1) Slightly paradoxically: if $\mathbb{P}(S_t) \geq \lambda$ then $\mathbb{E} \left(\frac{|V_t(N^a)|}{N} \right) \geq \frac{\lambda - a}{1 - a}$.

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By considering transpositions within \tilde{V}_t :

$$\mathbb{P}(I_t) = \frac{c' d}{N^d} \mathbb{E}(|\tilde{V}_t \cap W_t|) \geq \frac{c'}{N} (|\tilde{V}_t| - |W_t^c|)$$

